Self-Similar Relaxation Behavior at the Gel Point of a Blend of a Cross-Linking Poly(ε-caprolactone) Diol with a Poly(styrene-co-acrylonitrile)

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Received November 4, 1996; Revised Manuscript Received June 3, 1997

ABSTRACT: Novel polymer properties can be achieved by blending high molecular weight linear chains into a cross-linking system of short linear chains. This study is concerned with the rheological properties that are dominated at first by the highly entangled linear chains. However, with increasing extent of cross-linking, the short chains connect into a network structure and begin to dominate the rheology. The material here consists of cross-linking poly(ε-caprolactone) diol (PCL) and up to 40% of linear poly(styrene-co-acrylonitrile) (SAN) of high molecular weight. The blend was molecularly mixed before cross-linking. Three competing processes determine the structure of the system, (1) chemical cross-linking of the low molecular weight species into a sample spanning network of interpenetrating chains, (2) fluctuations in composition due to phase separation at increasing extents of reaction, and (3) crystallization of the PCL, which we tried to suppress as much as possible. At the gel point, systems with low SAN content show the typical scaling behavior of the critical gel with a self-similar relaxation spectrum, \( H(\lambda) = G(\lambda)n(\lambda/\lambda_0)^{-n}, \lambda > \lambda_0 \) at low probing frequencies, \( \omega < 1/\lambda_0 \). However, for the systems with high concentrations of the inert component, the self-similar region did not develop, possibly due to the phase separation induced by the cross-linking. The relaxation exponent, \( n \), decreased with increasing concentration of the highly entangled linear component. The results suggest that dynamic mechanical methods are applicable for determination of the gel point for homogeneous semi-IPN systems.

Introduction

Interpenetrating polymer networks (IPNs)\(^1\) have the potential for achieving a wide range of synergistic chemical and physical properties as required in specific polymer applications. IPNs are polymer blends in which the molecules of each component are cross-linked while being molecularly mixed with the other component(s). This is a unique way of conserving miscibility in temperature ranges in which the blend would otherwise phase separate. A semi-IPN is a special type of IPN comprising one cross-linkable polymer and a second nonreacting component of high molecular weight above its entanglement limit. Especially intriguing as novel materials are highly asymmetric semi-IPNs in which the inert polymer consists of highly entangled macro-molecules and the cross-linking molecules are very small initially. At the gel point (GP) of such a system one would expect good adhesion and damping properties from the developing network and good cohesion from the entangled chains. These properties could, for instance, be useful for an adhesive layer in composite materials. In another application, a small amount of high molecular weight polymer could be added to a low molecular weight cross-linking system (which by itself would cross-link into a brittle solid) with the purpose of decreasing the brittleness of the fully cross-linked material.

The preparation of an IPN is a challenging task because it involves two different transition phenomena that may interfere with each other, i.e., liquid-liquid phase separation between the components and sol-gel transition of at least one of the components. Recent experimental\(^2-4\) and theoretical\(^5,6\) studies specifically investigated the liquid-liquid phase separation process in semi-IPN systems. From these, it is widely understood that cross-linking may not only induce phase separation but also affect its kinetics. In order to control the final molecular structure, it is essential to understand the cross-linking behavior in greater detail.

Although numerous studies of IPNs were reported in the literature, most of them focused their efforts on the analysis of final properties and/or structures and less on the structure development. This study has the objective of determining the state of the sample at the sol-gel transition point, the so-called gel point, which is an important reference state for the development of new IPNs. The chemical gel point (GP) is defined as the instant at which the weight average molecular weight of the cross-linking component diverges to infinity. This may be detected chemically or physically. Beyond GP, the entire sample behaves as a solid even if only one of the components has solidified while the other component(s) is (are) still in the liquid state. In a multicomponent system such as an IPN, the presence of a second component might disturb the precise determination of GP in a yet unknown way.

Holly et al.\(^7\) developed a simple rheological method to determine GP. It is based on the observation that materials at GP (called "critical gels" [CG]) exhibit a self-similar relaxation time spectrum (CW-spectrum)\(^8-10\)

\[
H(\lambda) = \frac{S}{\Gamma(n)}\lambda^{-n}, \quad \text{for } \lambda_0 < \lambda < \infty \quad (1)
\]

which results in a relaxation modulus...
The gel stiffness, $S$, and the relaxation exponent, $n$, are the only material parameters necessary to completely characterize the linear viscoelastic properties at GP. $\lambda_0$ indicates the crossover from the universal self-similar behavior to some material specific short-time behavior.

The same scaling behavior is also apparent in dynamic mechanical experiments on critical gels where the storage modulus, $G'$, and the loss modulus, $G''$, are

$$G' = \frac{G''}{\tan \delta_c} = S \omega^n (1 - n) \cos \delta_c \quad 0 < \omega < 1/\lambda_0 \quad (3)$$

The phase angle between stress and strain, $\delta_c$, is independent of frequency, $\omega$, but proportional to the relaxation exponent $n$.

$$\delta_c = \pi n / 2 \quad (4)$$

Holly et al.\(^7\) probed the loss tangent of cross-linking polymers at a range of frequencies and identified GP as the instant at which eq 4 is satisfied, i.e., $\tan \delta$ curves intersect at GP.

Several studies explored this phenomenon and determined $S$ and $n$ values for single-component systems including poly(dimethylsiloxane) (PDMS),\(^8\) -11 polyurethanes,\(^12\) -16 epoxy systems,\(^17\) poly(propylene oxide),\(^18\) and even physically cross-linked polymers.\(^19\),\(^20\) The low-frequency self-similar behavior was observed in all these systems. The values of $n$ varied over the entire range between 0 and 1, depending on molecular compositions and cross-linking conditions. Short chains seem to give values of $n$ around 0.7, while long chains form critical gels with $n$ around 0.5.\(^21\)

Since the self-similar behavior of CG is universally observed in many different systems, it is expected to also occur in an IPN system. In order to apply Holly et al.'s method\(^7\) for determination of GP of the cross-linking component in an IPN system, however, we should know limitations of the self-similar behavior or any other effects caused by the presence of a second component in the IPN system.

Only few systematic studies were performed to investigate the effects of addition of a second component on the self-similar behavior of CG.\(^11\),\(^15\) Scanlan et al.\(^11\) explored divinyl-terminated PDMS systems of two different prepolymer molecular weights $M_w = 1 \times 10^4$ and $M_w = 4 \times 10^4$. Nonreacting PDMS with molecular weight of about $1.1 \times 10^4$ was added as an inert second component to each prepolymer. It was found that the critical self-similar behavior can be observed in the entire frequency window between 0.5 and 50 rad/s at any concentration of the inert component. However, $S$ and $n$ depend on the composition. For both prepolymer, $S$ monotonously decreases with increasing concentration of the inert component. $n$ is not affected by adding some of the inert component; however, high concentrations of the inert component result in very soft critical gels with an unusual high value of $n$, close to the theoretical limit of $n = 1$. Muller et al.\(^15\) at about the same time found that addition of an inert component of moderate $M_w$ increased the $n$ value in cross-linking poly(ethylene oxide) with isocyanate. In all cases, the systems seemed to be mixed on molecular level; i.e., phase separation phenomena were excluded from these experiments.

Gelation is somewhat more complicated for highly asymmetric semi-IPN systems in which the inert poly-
Superposed with a reference temperature of 140 °C, the dynamic storage modulus $G'$ and loss modulus $G''$ of SAN were measured at 140, 150, 160, 170, and 180 °C and superposed with a reference temperature of 140 °C. $G'$ ($\omega$) and $G''$ ($\delta$) of PCL were measured at 140 °C.

**Results**

**Properties of Premixtures.** The complex dynamic viscosity, $\eta^*$, of the PCL/SAN blend samples was measured at 110 °C with frequencies between 0.1 and 100 rad/s before addition of poly-IPDI, as shown in Figure 2. The onset of shear thinning, which is typical for high molecular weight polymers, was visible for the systems with high concentrations of SAN (FM60, FM70, FM80; the number next to FM denotes the weight percent of cross-linking component). The zero shear viscosity, $\eta_0$, was determined as the complex dynamic viscosity where it levels off at low frequencies.

DSC heating traces (Figure 3) suggested that PCL contained a small amount of crystal phase even after rapid quenching. The melting point of the crystal phase was 325 K, and the $T_g$ of the amorphous phase was 212 K. SAN is an amorphous polymer with a $T_g$ of 380 K. During heating, the PCL started to crystallize (DSC exotherm) and then it melted again (DSC endotherm). No significant transition was observed after the melting of the PCL crystal.

**Properties of Fully Cross-Linked Samples.** The dynamic complex moduli, $G^*$, of fully cross-linked samples with low SAN concentration (FM100, FM95, FM90 and FM80) were almost independent of frequency, as shown for 110 °C in Figure 4. For further evaluation of these systems, we adopted the value of $G^*$ measured at $\omega = 0.1$ rad/s as the equilibrium modulus, $G_e$. As far as FM70 and FM60, $G^*$ depended on frequency and $G_e$ could not be reduced to a single number.

The cross-linking reduces the crystallinity in the quenched samples. This can be seen when comparing the top curves in Figures 3 and 5. The addition of SAN further reduces the degree of crystallization in the PCL. This shows as endotherm in Figure 5. Up to an SAN concentration of 20 wt % (FM80) crystallinity can be detected in the quenched samples. For higher SAN content (FM70 and FM60) it seems to be suppressed completely.

The cross-linking of the PCL (FM100, FM95, FM90, FM80) increased the (single) glass transition temperature (see Figure 5). After cross-linking, FM70 and FM60 exhibited two glass transitions. $T_g$'s at low temperature did not increase significantly or were even lowered by the cross-linking. The higher $T_g$'s observed

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**Table 2. Composition of Blend Samples**

<table>
<thead>
<tr>
<th>Sample</th>
<th>FM100</th>
<th>FM95</th>
<th>FM90</th>
<th>FM80</th>
<th>FM70</th>
<th>FM60</th>
</tr>
</thead>
<tbody>
<tr>
<td>PCL/SAN/poly-IPDI (wt/wt/wt)</td>
<td>93.6/0/6.4</td>
<td>88.9/5/6.1</td>
<td>84.2/10/5.8</td>
<td>74.9/20/5.1</td>
<td>65.5/30/4.5</td>
<td>56.1/40/3.9</td>
</tr>
</tbody>
</table>

*The number next to FM denotes the weight percent of cross-linking component.
for FM70 and FM60 were 379 and 380 K, respectively. They were slightly lower than the $T_g$ of SAN, which was determined to be about 380 K.

Properties of Critical Gels. At the beginning of the rheological measurement, before cross-linking starts, the system has very low values of $G'$ and $G''$. With continuing cross-linking reaction, the values gradually increased and finally $G'$ exceeded $G''$. This is shown in Figure 6 for FM80. In the course, the tan $\delta$ decreases with reaction time, as shown in Figure 7 with the frequency as a parameter. The tan $\delta$ at low frequencies (between 0.1 and 10 rad/s) intersected at a certain time, $t_c$, which we will identify as gel point. For FM80, $t_c$ was found to be about 111 min. Although $t_c$ increased with increasing concentration of SAN, similar results were obtained for the other systems as long as the concentration of SAN was low enough (FM100, FM95, FM90, and FM70). However, such an instant at which tan $\delta$ became independent of frequency was not observed for FM60 in the accessible frequency window; i.e., the data gave no information about a possible GP.

Although $G'$ and $G''$ at $t_c$ seem to exhibit the self-similar behavior in about the entire experimental frequency range (see Figure 8), small deviations start to occur at about 10 rad/s for all the systems. This can be seen more readily on a plot of tan $\delta$ (Figure 9). An exception is FM100, which exhibited the self-similar behavior in the entire frequency range of the experiment.

$S$ and $n$ according to eqs 3 and 4 could be determined at $t_c$ for these systems containing 30 wt % of SAN or less, as summarized Table 3. The relaxation exponent took on a range of values between about 0.41 $\leq n \leq 0.61$. The corresponding gel stiffness values ranged from about $3.6 \times 10^2$ (Pa $^{0.61}$) to $1.1 \times 10^4$ (Pa $^{0.41}$). Figure 10 shows the relaxation exponent and the gel stiffness for the range of SAN content of our experiments. The exponent decreases with increasing SAN content, while
S increases with increasing SAN content for these samples except for FM70. The exponent of FM70 is between those of FM95 and FM90, and the gel stiffness is slightly smaller than that of FM80.

Discussion

Phase Separation Induced by Cross-Linking.

With continuing cross-linking reaction, the cross-linking component of the semi-IPN system increases its molecular weight and forms a network structure. The increase of the molecular weight reduces the gain of entropy of mixing. In addition, the network structure has to store the strain energy induced by swelling with a linear component. Since these additional energies make the homogeneous state unstable, molecularly mixed components before cross-linking might phase separate in the course of cross-linking. The increase of the molecular weight also changes the shape of the coexistence curve of the phase diagram. For linear chains of uniform molecular weights $M_A$ and $M_B$, the...
critical composition, $\Phi_c$, can be expressed as

$$\Phi_c = \frac{\sqrt{N_B}}{\sqrt{N_A} + \sqrt{N_B}}$$  \hspace{1cm} (5)$$

where $N_A$ and $N_B$ are the degree of polymerization of components A and B, respectively.\textsuperscript{24} Even if the above equation has been derived for linear molecules, it gives guidance for the system here where the chains of component A are highly branched and broadly distributed in size due to cross-linking. With an increase of the molecular weight of one of the components, the phase diagram becomes asymmetric. At the extreme of $N_A \to \infty$, $\Phi_c$ goes toward 0, suggesting that one of the phase-separated phases consists of its pure component.

The DSC results before cross-linking suggest that the PCL/SAN blends consist of crystalline PCL and an amorphous phase consisting of molecularly mixed PCL and SAN at temperatures below the crystal melting point of PCL. Since a single $T_g$ was observed for each un-cross-linked blend and the $T_g$ monotonically increased with increasing SAN content (see Figure 11), the system is supposed to be in its homogenous state at the cross-linking temperature of 110 °C.

After completion of the cross-linking reaction, FM70 and FM60 exhibit two $T_g$'s. As the asymmetric phase diagram of phase-separated semi-IPN suggests, the higher ones were almost equal to the $T_g$ of SAN itself and the change of heat capacity at $T_g$ was very small. The lower ones are slightly higher than the $T_g$ of cross-linked PCL. From these results, we may conclude that the systems phase separated into an almost pure SAN phase and a PCL rich phase containing some amount of SAN. However, we could not determine whether such transition from the homogeneous state to the partially phase separated state took place before or after GP of PCL. For cross-linked FM95, FM90, and FM80, there is no evidence for such phase separation.

**Detection of GP in a Semi-IPN System.** Theories of gelation\textsuperscript{25,26} predict a power law cluster number distribution at the gel point in terms of the cluster mass, $M$,

$$N_M \sim M^{-\gamma}$$  \hspace{1cm} (6)$$

where $\gamma$ is termed the polydispersity index. The cluster radius distribution is also expected to be a power law at GP

$$R_M \sim M^{1/d_f}$$  \hspace{1cm} (7)$$

where $d_f$ is the fractal dimension. The self-similar...
relaxation behavior of CG might originate from these statistical self-similarities. Several researchers \(^{13,17,27}\) derived expressions of the relaxation exponent in terms of the fractal dimension by assuming that \(N_M\) and \(M\) independently relate to the spectrum \(H(\lambda)\). A critical gel consisting of only cross-linkable component satisfies those eqs over a very wide range of mass. This leads to a very wide time range where the self-similar relaxation behavior can be observed.

If we add to CG a linear polymer that does not satisfy eq 6, the power law cluster number distribution can be observed only at high molecular weights, above the molecular weight of the linear component. This explains the restricted time window of the self-similar relaxation behavior in a semi-IPN system containing long linear molecules.

In fact, we found that the lower limiting relaxation time, \(\lambda_0\), has values in the experimental frequency range for FML95, FM90, FM80, and FM70. In comparison, the \(\lambda_0\) of FM100, which does not contain SAN, is outside of the accessible frequency range. We should, however, emphasize here that the self-similar relaxation behavior can be observed at low frequencies for these systems. This suggests that these systems at their gel point still possess the typical self-similarity if we probe spatial sizes larger than the size of the SAN molecule. Therefore, the observed \(t_c\) is considered to be the gel time of PCL.

For the premixtures, shear thinning was observed only for FM80, FM70, and FM60 and occurred at particular frequencies, depending on the composition (see Figure 2). However, \(\omega_0 = 1/\lambda_0\) is observed even for FM95 and FM90 and does not seem to depend on the composition very much (see Figure 9). This might be explained as follows: The transition of relaxation behavior from entanglement to flow region is an expression of the longest relaxation time of SAN. This depends on the composition since PCL acts as a diluent of SAN in the premixtures.

The DSC results suggest that phase separation induced by cross-linking takes place in FM70 and FM60. Only for FM60, however, could the self-similar relaxation behavior not be observed. This suggests that the phase separation may already start before the FM60 reaches GP. If the phase separation takes place before GP, we would not expect the self-similar behavior within our experimental frequency window. The phase separation might affect the self-similar behavior of FM70 as well. This would explain its peculiar set of \(n\) and \(S\) values that deviate from the tendency observed on the samples with low SAN concentration. A careful investigation of the above hypothesis would require measurements of the structure growth during cross-linking, possibly by light scattering.

In the next two sections, we will discuss the values of \(S\) and \(n\) that were determined for the system with low concentrations of SAN (FM100, FM95, FM90, and FM80).

**Effect of Addition of Linear Polymer on the Relaxation Exponent.** This study shows remarkable contrast to Scanlan’s observation who found that \(S\) decreases and \(n\) increases with increasing concentration of the nonreacting linear polymer. This might be attributed to the difference of molecular weights of the additives. In the case of Scanlan’s PDMS system, the molecular weight of the un-cross-linkable PDMS is below the entanglement molecular weight of about 24 500. \(^{28}\) Therefore, the PDMS acts as a diluent in the system. On the other hand, the molecular weight of SAN in this study is much above the entanglement molecular weight, which was reported to be 11 000 for SAN with 24.1 wt % of acrylonitrile. \(^{29}\) Therefore, the dilution effect is masked by the entanglement behavior of the high molecular weight SAN. Molecular weight of a linear component might be an important factor that determines whether \(n\) increases or decreases with increasing concentration of linear component.

Durand et al.\(^ {13}\) and Martin et al.\(^ {17}\) predicted \(n\) in terms of the fractal dimension as

\[
\eta = \frac{d}{d + 2} \quad (8)
\]

where \(d\) is the space dimension (=3). Muthukumar\(^ {27}\) considered the excluded volume effect and suggested that, if the excluded volume effect is fully screened, the relaxation exponent becomes

\[
\eta = \frac{d(d + 2 - 2d_f)}{2(d + 2 - d_f)} \quad (9)
\]

These equations suggest that the relaxation exponent increases with decreasing fractal dimension. Our observation, therefore, can be qualitatively elucidated: As the SAN concentration is increased, the structure becomes more dense and \(d_f\) is higher, leading to a decrease in \(n\).

**Relationships between \(S\) and \(n\).** Many observations indicated that \(S\) and \(n\) are coupled. Because of its physical dimension of Ps \(s^n\), it was apparent that the stiffness \(S\) of the critical gel was composed of a modulus and a time constant. \(^{8}\) An experimental relationship between \(S\) and \(n\) was found for cross-linking PDMS systems\(^ {11}\) as

\[
S = G_{cg} \lambda_{cg}^n \quad (10)
\]

where the meaning of the modulus \(G_{cg}\) and the time constant \(\lambda_{cg}\) (subscript cg for critical gel) needs to be specified. It was not completely clear how this relation can be extended to semi-IPNs or whether it is applicable at all. This was examined by plotting \(\log S\) against \(n\) (see Figure 12). In fact, we can observe fairly good linearity between \(n\) and \(\log S\). A least-squares fit (a solid line) would give \(G_{cg} = 6.31 \times 10^6\) Pa and \(\lambda_{cg} = 1.90 \times 10^{-8}\) s.

It was found that \(G_{cg}\) was very close to the modulus of the fully cross-linked material, \(G_e\). \(^ {11}\) The most simple...
assumptions for $\lambda_{cg}$ were that it is either the low-frequency end of the power law, as defined in eqs 1–3, the longest relaxation time of the precursor, or the ratio of precursor viscosity and modulus, $\eta_p/G_e$.11 The vastly different relaxation time scales of the two components in our system allow us to further explore the meaning of $\lambda_{cg}$.

Recently, the relationship $\lambda_{cg} = \eta_p/G_e$ was investigated16 in detail for PCL systems with different molecular weight precursors. It was found that the relationship between $S$ and $n$ can be represented approximately in terms of $G_e$ and $\eta_p$ as

$$S = G_e(\eta_p/G_e)^n$$  \hspace{1cm} (11)

where the two parameters are functions of molecular weight. $S$ values according to eq 11 were calculated by using experimentally determined values of $n$, $G_e$, and $\eta_p$ for each system (see Figure 13). The calculated values (crosses on broken lines) are shown in comparison with the corresponding experimental $S$ values (squares) in Figure 12, as well. There are only small discrepancies between the calculated values and the experimental values. We may conclude that eq 11 is a reasonably good expression for the relationship between $S$ and $n$, which is applicable even for the systems containing a long linear component.

Scanlan et al.'s11 other two assumptions for $\lambda_{cg}$ do not seem to apply to these semi-IPN systems, where the crossover time $\lambda_0 (= 1/\omega_0)$ is of the order of $10^{-1}$ (see Figure 9). The longest relaxation time of the precursor, $\lambda_p$, is of the same order as $\lambda_0$ and, hence, is also much larger than $\lambda_{cg}$. In such a case, it is concluded that eq 10 cannot predict the value $\lambda_p$.

Conclusions

These results suggest that the gel point of a crosslinking component in a semi-IPN system is detectable by its self-similar relaxation behavior, as long as the system is homogeneous. From a practical point of view, this method is advantageous because it is simpler than conventional methods and allows simultaneous measurement during processing.

The self-similar relaxation behavior is observed in cross-linking poly(ε-caprolactone) and poly(styrene-co-acrylonitrile) blends with low concentrations of the linear component from the restricted time window. Phase separation might hinder the self-similar behavior for the systems with a high concentration of the linear component.

The relaxation properties, $n$ and $S$, strongly depend on composition. The relaxation exponent decreases with increasing concentration of the linear component, suggesting that the fractal dimension increases with increasing concentration of the linear component. The molecular weight of the nonreacting linear component might be an important factor that determines whether $n$ increases or decreases with increasing concentration of the linear component.

An experimental relationship between $n$ and $S$, $S = G_e(\eta_p/G_e)^n$, is observed in this system as was observed in 100% poly(ε-caprolactone) critical gels. Here, $\eta_p$ is the zero shear viscosity of the mixture at the onset of the cross-linking reaction and $G_e$ is the final modulus.

Acknowledgment. We gratefully acknowledge Daicel Chemical Industries, Ltd. for financial support and providing samples. We also thank Huls AG for samples of the cross-linker.

References and Notes

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